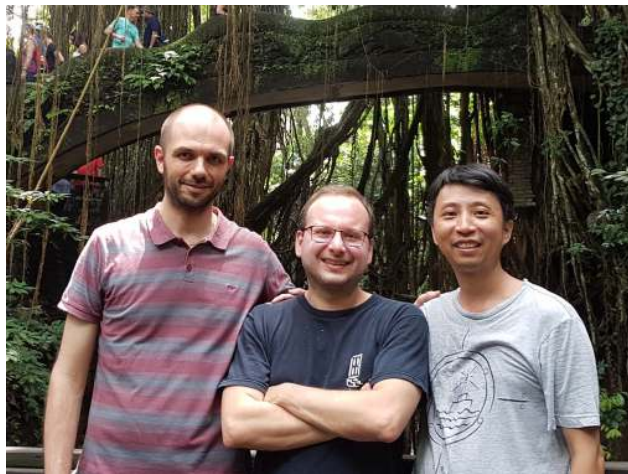


The critical 2d
Stochastic Heat Flow
(is not a G.M.C.)

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Based on joint works with
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OVERVIEW

I. THE CRITICAL 2D STOCHASTIC HEAT FLOW

Motivations, construction, Key features

II. MAIN RESULTS

Gaussian Multiplicative Chaos, moment bounds

III. IDEAS FROM THE PROOF

REFERENCES

- [CSZ 21] F. Caravenna, R. Sun, N. Zygouras
THE CRITICAL 2D STOCHASTIC HEAT FLOW
arXiv (2021)
- [CSZ 22] F. Caravenna, R. Sun, N. Zygouras
... IS NOT A GAUSSIAN MULTIPLICATIVE CHAOS
In preparation

1. WHAT IS THE STOCHASTIC HEAT FLOW ?

It is a "universal" stochastic process of measures on \mathbb{R}^2

$$\mathbb{Z}_{s,t}^g(dy, dx) \quad 0 \leq s < t < \infty \quad x, y \in \mathbb{R}^2 \quad g \in \mathbb{R}$$

Heuristics: "natural candidate" solution of the (ill-defined)
"critical" 2d Stochastic Heat Equation:

$$\partial_t U(t, x) = \frac{1}{4} \Delta U(t, x) + \beta \underbrace{\xi(t, x)}_{\text{SPACE-TIME WHITE NOISE}} U(t, x) \quad (\text{SHE})$$

Fix starting time s , initial condition $U(s, x) = \psi(x)$:

$$\rightsquigarrow \int_{\mathbb{R}^2} \psi(y) \mathbb{Z}_{s,t}^g(dy, \cdot) = \text{"natural candidate" for } U(t, \cdot)$$

($U(t, \cdot)$ expected not to be a function, but a measure on \mathbb{R}^2)

What do we mean by "natural candidate"? We will see by
REGULARIZATION of the eq. & RENORMALIZATION of $\beta \leftrightarrow g$

Henceforth we fix $s=0$, $\psi \equiv 1$ and call

$$\int_{y \in \mathbb{R}^2} \mathbb{Z}_{s,t}^g(dy, dx) = \mathbb{Z}_t^g(dx) = \text{STOCHASTIC HEAT FLOW (SHF)}$$

2. A LINK WITH DIRECTED POLYMERS

We can **REGULARIZE** the Stochastic Heat Equation by **discretizing time and space** (or by mollification \rightsquigarrow later)

$$\underbrace{\partial_t^N U(t,x)}_{N \cdot \{U(t,x) - U(t-\frac{1}{N}, x)\}} = \underbrace{\frac{1}{4} \Delta^N U(t,x)}_{\text{SCALED LATTICE LAPLACIAN}} + \underbrace{\beta N \langle \overbrace{\eta(t,x)}^{\text{I.I.D. ZERO MEAN, UNIT VARIANCE}} U(t,x) \rangle}_{\text{SPACE AVERAGE } \frac{1}{4} \sum_{x' \sim x}}$$

$$\text{for } (t,x) = \left(\frac{n}{N}, \frac{z}{\sqrt{N}} \right) \in \frac{\mathbb{N}}{N} \times \frac{\mathbb{Z}^2}{\sqrt{N}}$$

$$\text{Special choice of noise: } \eta(t,x) = \frac{e^{\overbrace{\beta \omega(n,z)}^{\text{I.I.D. } N(0,1)} - \frac{\beta^2}{2}} - 1}{\beta}$$

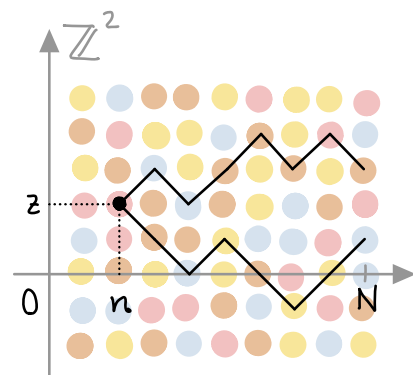
actually $1-t$

$$\uparrow U(t,x) = Z_N(n,z) = E \left[e^{\sum_{i=n+1}^N \beta \omega(i, S_i) - \frac{\beta^2}{2}} \mid S_n = z \right]$$

(S_i) SIMPLE RANDOM WALK ON \mathbb{Z}^2

Partition function of the

DIRECTED POLYMER
IN RANDOM ENVIRONMENT



3. MAIN RESULTS: EXISTENCE OF THE STOCHASTIC HEAT FLOW

The solution of the discretized Stochastic Heat Equation is

$$u(t, x) = Z_N(\underbrace{tN}_{\text{INTEGER PART}}, \underbrace{x\sqrt{N}}_{\text{INTEGER PART}}) \quad (t, x) = \left(\frac{t}{N}, \frac{x}{\sqrt{N}}\right)$$

A "natural candidate" solution of the original equation (ill-defined) is obtained removing the discretization: do we have

$$\int \varphi(x) Z_N(tN, x\sqrt{N}) dx \xrightarrow[N \rightarrow \infty]{d} ?$$

Not yet! As we remove the discretization, we need to RENORMALIZE the noise strength β (intermediate disorder regime):

$$\beta \sim \hat{\beta} \frac{\sqrt{\pi}}{\sqrt{\log N}} \quad \text{for fixed } 0 < \hat{\beta} < \infty$$

Indeed, for $x \neq y \in \mathbb{R}^2$

$$\mathbb{E}[Z_N(tN, x\sqrt{N}), Z_N(tN, y\sqrt{N})] \xrightarrow[N \rightarrow \infty]{} \begin{cases} 0 & \text{if } \hat{\beta} < 1 \\ K(x, y) \in (0, \infty) & \text{if } \hat{\beta} = 1 \\ \infty & \text{if } \hat{\beta} > 1 \end{cases}$$

$$(\mathbb{E}[Z_N(tN, x\sqrt{N})] \equiv 1)$$

[Bertini-Cancrini '98]

The 2d directed polymer model undergoes a phase transition from weak to strong disorder. (see [CSZ, AAP 2017])

We fix the critical point $\hat{\beta} = 1$, in fact a critical window:

$$\textcircled{\star} \quad \beta = \left(1 + \frac{\mathcal{G} + o(1)}{\log N} \right) \frac{\sqrt{\pi}}{\sqrt{\log N}} \quad \text{for } -\infty < \mathcal{G} < +\infty$$

Our first main result:

Theorem 1. [CSZ 21]

Fix $\mathcal{G} \in \mathbb{R}$. As $N \rightarrow \infty$ we have the convergence in f.d.d. of

$$\left(Z_N(tN, x\sqrt{N}) dx \right)_{t \in [0,1]}$$

(as random measures on \mathbb{R}^2) to a unique limit

$$\left(\mathcal{Z}_t(dx) \right)_{t \in [0,1]}$$

which we call the CRITICAL 2D STOCHASTIC HEAT FLOW.

The SHF encodes asymptotic properties of directed polymers, and it gives a meaning to the solution of the Stochastic Heat Equation.

The rest of the talk is devoted to investigating the SHF.

4. GAUSSIAN MULTIPLICATIVE CHAOS

Let $X = (X(x))_{x \in \mathbb{R}^2}$ be a (generalized) Gaussian field on \mathbb{R}^2 .
Fix a reference measure $\mu(dx)$ on \mathbb{R}^2 .

Gaussian Multiplicative Chaos (GMC) is a random measure on \mathbb{R}^2 :

$$\text{formally } \mathcal{M}(dx) = e^{X(x) - \frac{1}{2} \text{VAR}[X(x)]} \mu(dx)$$

Often X is specified by (cf. GFF)

$$\text{Cov}[X(x), X(y)] =: \kappa(x, y) \rightarrow \infty \quad \text{as } |y-x| \rightarrow 0$$

$X(x)$ is not defined pointwise: $X(\varphi) = \int \varphi(x) X(x) dx \sim N(0, \sigma_\varphi^2)$

$$\sigma_\varphi^2 = \text{VAR}[X(\varphi)] = \iint \varphi(x) \varphi(y) \kappa(x, y) dx dy < \infty$$

Approximate $\kappa(x, y) = \lim_{\varepsilon \downarrow 0} \kappa_\varepsilon(x, y)$ by smooth covariance functions.

Then $X_\varepsilon \sim N(0, \kappa_\varepsilon)$ is defined pointwise, $X_\varepsilon(x)$, and we set

$$\mathcal{M}(dx) := \lim_{\varepsilon \downarrow 0} e^{X_\varepsilon(x) - \frac{1}{2} \text{VAR}[X_\varepsilon(x)]} \mu(dx)$$

5. MAIN RESULTS: SHF & GMC

SUB-CRITICAL REGIME

For $\hat{\beta} < 1$, it is known that $\log Z_N$ has Gaussian fluctuations:

EDWARDS-WILKINSON

$$\int \varphi(x) \sqrt{\log N} \left(\log Z_N(t_N, x\sqrt{N}) - \mathbb{E}[\log Z_N] \right) dx \xrightarrow[N \rightarrow \infty]{d} \text{GAUSSIAN}$$

This can be roughly rephrased as:

$$Z_N(t_N, x\sqrt{N}) \approx \exp(\text{GAUSSIAN})$$

Does this picture apply at criticality, i.e. for $\hat{\beta} = 1$?

We can state our second main result:

Theorem 2. [CSZ 22]

The critical 2d SHF is **not** a GMC.

KPZ EQUATION

This suggests that $\log Z_N(t_N, x\sqrt{N})$ is **not** asymptotically Gaussian at criticality $\hat{\beta} = 1$.

Also: the SHF $\mathcal{Z}_t^\nu(dx)$ is a new class of random measures on \mathbb{R}^2 .

6. MAIN RESULTS: MOMENT BOUNDS

Consider a GMC $\mathcal{M}(dx) = e^{X(x) - \frac{1}{2} \text{VAR}[X(x)]} \mu(dx)$.
COVARIANCE $K(x,y)$

We can compute the moments

$$\mathbb{E}[\mathcal{M}(dx)] = \mu(dx)$$

$$\mathbb{E}[\mathcal{M}(dx) \mathcal{M}(dy)] = e^{K(x,y)} \mu(dx) \mu(dy)$$

Henceforth we fix $t > 0$, $g \in \mathbb{R}$. For the SHF $\mathcal{Z}_t^g(dx)$ we have

$$\mathbb{E}[\mathcal{Z}_t^g(dx)] = dx$$

$$\mathbb{E}[\mathcal{Z}_t^g(dx) \mathcal{Z}_t^g(dy)] = K(x,y) dx dy$$

$$\lim_{N \rightarrow \infty} \mathbb{E}[Z_N(tN, x\sqrt{N}), Z_N(tN, y\sqrt{N})] \sim C \log \frac{1}{|y-x|} \text{ as } |y-x| \rightarrow 0$$

We can match 1st & 2nd moments of \mathcal{M} and \mathcal{Z}_t^g by choosing

$$\mu(dx) = dx \quad K(x,y) = \log \log \frac{1}{|y-x|}$$

The GMC \mathcal{M} is now completely determined.

Set $\mathcal{L}_t^\vartheta(\varphi) := \int \varphi(x) \mathcal{L}_t^\vartheta(dx)$, $\mathcal{M}(\varphi) := \int \varphi(x) \mathcal{M}(dx)$

We can prove that $\mathcal{L}_t^\vartheta \neq \mathcal{M}$ by comparing **third moments**.

Theorem 3. (3rd moment bound)

[CSZ 22]

If φ is the indicator of a ball, or the heat kernel, then

$$\mathbb{E}[\mathcal{L}_t^\vartheta(\varphi)^3] > \mathbb{E}[\mathcal{M}(\varphi)^3].$$

For the heat kernel $g_\delta(x) := \frac{e^{-\frac{|x|^2}{2\delta}}}{2\pi\delta}$ we bound higher moments.

Theorem 4. (higher moments bound)

[CSZ 22]

There is $\eta = \eta_{t,\vartheta} > 0$ such that $\forall h \in \mathbb{N}$

$$\mathbb{E}[\mathcal{L}_t^\vartheta(g_\delta)^h] \geq (1+\eta) \mathbb{E}[\mathcal{L}_t^\vartheta(g_\delta)^2]^{\binom{h}{2}} \quad \forall \delta \in (0,1),$$

$$\text{while} \quad \mathbb{E}[\mathcal{M}(g_\delta)^h] \sim \mathbb{E}[\mathcal{M}(g_\delta)^2]^{\binom{h}{2}} \quad \text{as } \delta \downarrow 0.$$

Moments **upper** bounds for $\beta < 1$ by [Cosco, Zeitouni], [Lygkonis, Zygouras]

7. IDEAS FROM THE PROOF (1): MOMENT FORMULAS

We prove Theorem 3 exploiting an **exact formula** for the third moment, derived in [CSZ 19]:

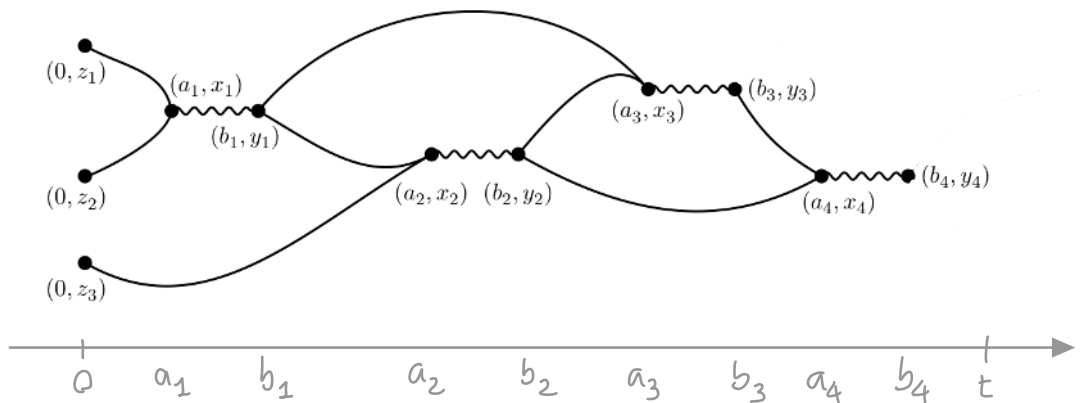
$$\mathbb{E} \left[\mathcal{L}_t^\vartheta(dx) \mathcal{L}_t^\vartheta(dy) \mathcal{L}_t^\vartheta(dz) \right] = K^{(3)}(x, y, z) dx dy dz$$

$$K^{(3)}(z_1, z_2, z_3) = \sum_{m \geq 2} 3 \cdot 2^{m-1} (2\pi)^m \int \dots \int d\vec{a} d\vec{b} d\vec{x} d\vec{y} g_{\vec{z}}^{(m)}(\vec{a}, \vec{b}, \vec{x}, \vec{y})$$

$$0 < a_1 < b_1 < \dots < a_m < b_m < t$$

$$x_1, y_1, \dots, x_m, y_m \in \mathbb{R}^2$$

$g^{(4)} =$



where:

$$g_{\frac{a-b}{2}}(x-y)$$

$$g_{\frac{b-a}{4}}(y-x) \cdot G_\vartheta(b-a)$$

By an explicit Gaussian integration, we can integrate out all space variables \vec{x}, \vec{y} to get the equality

$$\int \dots \int g_{\vec{z}}^{(m)}(\vec{a}, \vec{b}, \vec{x}, \vec{y}) d\vec{x} d\vec{y} = \prod_{i=1}^m G_g(b_i - a_i) \cdot g_{a_1}(z_1 - z_2) \cdot g_{\bar{a}_2}\left(z_3 - \frac{z_1 + z_2}{2}\right) \cdot \prod_{i=3}^m g\left(\frac{0}{\overline{a_i - b_{i-2}}}\right)$$

for suitable $\bar{a}_2 < a_2$ and $\overline{a_i - b_{i-2}} < a_i - b_{i-2}$.

Exploiting monotonicity of the heat Kernel, we replace the quantities \bar{a}_2 by a_2 and $\overline{a_i - b_{i-2}}$ by $a_i - b_{i-2}$.

This yields the bound $\mathbb{E}[\mathcal{L}_t^\vartheta(\varphi)^3] > \mathbb{E}[\mathcal{M}(\varphi)^3]$. \square

Remark. Exact formulas are available for higher moments:

$$\mathbb{E}\left[\prod_{i=1}^h \mathcal{L}_t^\vartheta(dz_i)\right] = K^{(h)}(z_1, \dots, z_h) dz_1 \dots dz_h$$

$$K^{(h)}(z_1, \dots, z_h) = \sum_m \text{EXPLICIT (COMPLICATED) INTEGRALS ON } (\text{TIME} \times \text{SPACE}^h)^{2m}$$

This was proved by [Gu, Quastel, Tsai, PMP'21]

8. IDEAS FROM THE PROOF (II): CORRELATION INEQUALITY

Recall that

$$Z_N(\phi, z) = E \left[e^{\sum_{\ell=1}^N \beta \omega(\ell, S_\ell) - \frac{\beta^2}{2}} \mid S_0 = z \right]$$

Then

$$\boxed{\star} \quad E \left[\prod_{i=1}^h Z_N(z_i) \right] = E \left[\prod_{1 \leq i < j \leq h} e^{\beta^2 L_N^{(i,j)}} \mid S_0^{(i)} = z_i \quad \forall i \right]$$

$$L_N^{(i,j)} = \sum_{\ell=1}^N \mathbb{1}_{\{S_\ell^{(i)} = S_\ell^{(j)}\}} \quad \text{OVERLAP OF } S^{(i)} \text{ AND } S^{(j)}$$

$$\text{Our goal is} \quad E \left[\mathcal{Z}_t^\phi(g_\delta)^h \right] \geq (1+\eta) E \left[\mathcal{Z}_t^\phi(g_\delta)^2 \right]^{\frac{h}{2}}$$

Set $\eta=0$ (for simplicity) and replace \mathcal{Z}_t^ϕ by Z_N .

We need to exchange E and $\prod_{1 \leq i < j \leq h}$ in $\boxed{\star}$:

are $\exp(\beta^2 L_N^{(i,j)})$ positively correlated? (at least as $N \rightarrow \infty$)

We prove that this holds in a continuum setting, replacing

random walks $S^{(i)}$ \rightsquigarrow Brownian motions $B^{(i)}$

Based on the GAUSSIAN CORRELATION INEQUALITY [Rogers '14].

CONCLUSIONS

We introduced the CRITICAL 2D STOCHASTIC HEAT FLOW as a **scaling limit** of directed polymer partition Functions

It is a universal process of random measures on \mathbb{R}^2 , different from GMC, which is a natural candidate for the solution of the critical 2d Stochastic Heat Equation.

It has some explicit (but challenging!) features.

Many interesting questions are open:

- SINGULARITY W.R.T. LEBESGUE MEASURE
- FLOW PROPERTY
- CHARACTERIZING PROPERTIES
- TAKING LOG \rightsquigarrow KPZ

Thanks!