

Polynomial Chaos and Scaling Limits of Disordered Systems

2. Continuum model and free energy estimates

Francesco Caravenna

Università degli Studi di Milano-Bicocca

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Overview

In the previous lecture we saw how to build **continuum partition functions**

$$\mathcal{Z}_\delta^\omega \xrightarrow[\delta \rightarrow 0]{d} \mathcal{Z}^W \text{ (scaling limits of discrete partition functions)}$$

In this lecture we present two interesting applications of \mathcal{Z}^W

- ▶ Scaling limit of the full probability measure $\mathbf{P}_\delta^\omega \xrightarrow[\delta \rightarrow 0]{d} \mathcal{P}^W$
constructing a continuum version of the disordered system

We will focus on the DPRE [Alberts, Khanin, Quastel 2014b]
drawing inspiration from the Pinning [C., Sun, Zygouras 2016]

- ▶ Sharp asymptotics on the discrete model, in terms of **free energy** and **critical curve**

For this we will focus on **Pinning models** (rather than DPRE)

Outline

1. White noise and Wiener chaos
2. Continuum partition functions
3. The continuum DPRE
4. Pinning models

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White noise (1 dim.)

We are familiar with (1-dim.) **Brownian motion** $B = (B(t))_{t \geq 0}$

We are interested in its derivative “ $W(t) := \frac{d}{dt}B(t)$ ” called **white noise**

Think of W as a stochastic process $W = (W(\cdot))$ indexed by

$$\text{Intervals } I = [a, b] \quad \longmapsto \quad W(I) = B(b) - B(a) \sim \mathcal{N}(0, b - a)$$

$$\text{Borel sets } A \in \mathcal{B}(\mathbb{R}) \quad \longmapsto \quad W(A) = \int_{\mathbb{R}} \mathbb{1}_A(t) dB(t) \sim \mathcal{N}(0, |A|)$$

W is a **Gaussian** process with

$$\mathbb{E}[W(A)] = 0 \quad \text{Cov}[W(A), W(B)] = |A \cap B|$$

White noise

White noise on \mathbb{R}^d

It is a Gaussian process $W = (W(A))_{A \in \mathcal{B}(\mathbb{R}^d)}$ with

$$\mathbb{E}[W(A)] = 0 \quad \text{Cov}[W(A), W(B)] = |A \cap B|$$

$$\blacktriangleright \forall (A_n)_{n \in \mathbb{N}} \text{ disjoint} \implies W\left(\bigcup_{n \in \mathbb{N}} A_n\right) \stackrel{\text{a.s.}}{=} \sum_{n \in \mathbb{N}} W(A_n)$$

Almost a random signed measure on \mathbb{R}^d ... (but not quite!)

We can define single stochastic integrals $W(f) := \int f(x) W(dx)$

$$\mathbb{E}[W(f)] = 0 \quad \mathbb{E}[W(f)^2] = \|f\|_{L^2(\mathbb{R}^d)}^2$$

Multiple stochastic integrals

We can define

$$W^{\otimes k}(g) = \int_{(\mathbb{R}^d)^k} g(x_1, \dots, x_k) W(dx_1) \cdots W(dx_k)$$

For $d = 1$ we can restrict $x_1 < x_2 < \dots < x_k \rightsquigarrow$ iterated Ito integrals

For **symmetric** functions we have

$$\mathbb{E}[W^{\otimes k}(g)] = 0 \quad \mathbb{E}[W^{\otimes k}(g)^2] = k! \|g\|_{L^2((\mathbb{R}^d)^k)}^2$$

$$\text{Cov}[W^{\otimes k}(f), W^{\otimes k'}(g)] = 0 \quad \forall k \neq k'$$

Wiener chaos expansion

Any r.v. $X \in L^2(\Omega_W)$ measurable w.r.t. $\sigma(W)$ can be written as

$$X = \sum_{k=0}^{\infty} \frac{1}{k!} W^{\otimes k}(f_k) \quad \text{with} \quad f_k \in L_{\text{sym}}^2((\mathbb{R}^d)^k)$$

Discrete sums and stochastic integrals

Consider a lattice $\mathbb{T}_\delta \subseteq \mathbb{R}^d$ whose cells have volume $v_\delta \rightarrow 0$

Take i.i.d. random variables $(X_z)_{z \in \mathbb{T}_\delta}$ with zero mean and unit variance

Consider the “stochastic Riemann sum” (multi-linear polynomial)

$$\Psi_\delta := \sum_{\substack{(z_1, \dots, z_k) \in (\mathbb{T}_\delta)^k \\ z_i \neq z_j \ \forall i \neq j}} f(z_1, \dots, z_k) X_{z_1} X_{z_2} \cdots X_{z_k}$$

where $f \in L^2(\mathbb{R}^d)$ is (say) continuous.

$$(\sqrt{v_\delta})^k \Psi_\delta \xrightarrow[\delta \rightarrow 0]{d} \int_{(\mathbb{R}^d)^k} g(z_1, \dots, z_k) W(dz_1) \cdots W(dz_k)$$

(Check the variance!)

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Continuum partition function for DPRE

1d rescaled RW $S_t^\delta := \sqrt{\delta} S_{t/\delta}$ lives on $\mathbb{T}_\delta = ([0, 1] \cap \delta\mathbb{N}_0) \times \sqrt{\delta}\mathbb{Z}$

$$\begin{aligned} \mathbf{Z}_\delta^\omega &= \mathbf{E}^{\text{ref}} \left[\exp \left(\mathcal{H}^\omega \right) \right] = \mathbf{E}^{\text{ref}} \left[\exp \left(\sum_{n=1}^N (\beta \omega_{(n, S_n)} - \lambda(\beta)) \right) \right] \\ &= 1 + \sum_{(t, x) \in \mathbb{T}_\delta} \mathbf{P}^{\text{ref}}(S_t^\delta = x) \mathbf{X}_{t, x} \\ &\quad + \frac{1}{2} \sum_{(t, x) \neq (t', x') \in \mathbb{T}_\delta} \mathbf{P}^{\text{ref}}(S_t^\delta = x, S_{t'}^\delta = x') \mathbf{X}_{t, x} \mathbf{X}_{t', x'} + \dots \end{aligned}$$

Recall the LLT: $\mathbf{P}^{\text{ref}}(S_n = x) \sim \frac{1}{\sqrt{n}} g\left(\frac{x}{\sqrt{n}}\right)$ with $g(z) = \frac{e^{-\frac{z^2}{2}}}{\sqrt{2\pi}}$

$$\mathbf{P}^{\text{ref}}(S_t^\delta = x) = \mathbf{P}^{\text{ref}}(S_{\frac{t}{\delta}} = \frac{x}{\sqrt{\delta}}) \sim \sqrt{\delta} g_t(x) \quad g_t(x) = \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}}$$

Replacing $\mathbf{X}_{t, x} = e^{(\beta \omega_{(t, x)} - \lambda(\beta))} - 1 \approx \beta \mathbf{Y}_{t, x}$ with $\mathbf{Y}_{t, x}$ i.i.d. $\mathcal{N}(0, 1)$

Continuum partition function for DPRE

$$\begin{aligned} Z_N^\omega = & 1 + \beta\sqrt{\delta} \sum_{(t,x) \in \mathbb{T}_\delta} g_t(x) Y_{t,x} \\ & + \frac{1}{2} (\beta\sqrt{\delta})^2 \sum_{(t,x) \neq (t',x') \in \mathbb{T}_\delta} g_t(x) g_{t'-t}(x' - x) Y_{t,x} Y_{t',x'} + \dots \end{aligned}$$

Cells in \mathbb{T}_δ have volume $v_\delta = \delta\sqrt{\delta} = \delta^{\frac{3}{2}} \rightsquigarrow$ “Stochastic Riemann sums”
 converge to stochastic integrals if $\beta\sqrt{\delta} \approx \sqrt{v_\delta}$

$$\boxed{\beta \sim \hat{\beta} \delta^{\frac{1}{4}} = \frac{\hat{\beta}}{N^{\frac{1}{4}}}}$$

$$\begin{aligned} Z_N^\omega & \xrightarrow[\delta \rightarrow 0]{d} Z^W = 1 + \hat{\beta} \int_{[0,1] \times \mathbb{R}} g_t(x) W(dtdx) \\ & + \frac{\hat{\beta}^2}{2} \int_{([0,1] \times \mathbb{R})^2} g_t(x) g_{t'-t}(x' - x) W(dtdx) W(dt'dx') \\ & + \dots \end{aligned}$$

Constrained partition functions

We have constructed \mathcal{Z}^W = “free” partition function on $[0, 1] \times \mathbb{R}$
RW paths starting at $(0, 0)$ with no constraint on right endpoint

$$\mathcal{Z}^W = \mathcal{Z}^W((0, 0), (1, \star)) \quad \mathbb{E}[\mathcal{Z}^W] = 1$$

Consider now **constrained** partition functions: for $(s, y), (t, x) \in [0, 1] \times \mathbb{R}$

$$\text{Discrete:} \quad \mathcal{Z}_\delta^w((s, y), (t, x)) = \mathbf{E}^{\text{ref}} \left[\exp \left(\mathcal{H}^w \right) \mathbb{1}_{\{S_t^\delta = x\}} \middle| S_s^\delta = y \right]$$

Divided by $\sqrt{\delta}$, they converge to a continuum limit:

$$\mathcal{Z}^W((s, y), (t, x)) \quad \mathbb{E}[\mathcal{Z}^W((s, y), (t, x))] = g_{t-s}(x - y)$$

This is a function of white noise in the stripe $W([s, t] \times \mathbb{R})$

Four-parameter random process $\mathcal{Z}^W((s, y), (t, x)) \rightsquigarrow$ regularity?

Key properties

Key properties

For a.e. realization of W the following properties hold:

- ▶ **Continuity:** $\mathcal{Z}^W((s, y), (t, x))$ is jointly continuous in (s, y, t, x) (on the domain $s < t$)
- ▶ **Positivity:** $\mathcal{Z}^W((s, y), (t, x)) > 0$ for all (s, y, t, x) satisfying $s < t$
- ▶ **Semigroup** (Chapman-Kolmogorov): for all $s < r < t$ and $x, y \in \mathbb{R}$

$$\mathcal{Z}^W((s, y), (t, x)) = \int_{\mathbb{R}} \mathcal{Z}^W((s, y), (r, z)) \mathcal{Z}^W((r, z), (t, x)) dz$$

(Inherited from discrete partition functions: [drawing!](#))

How to prove these properties?

The 1d Stochastic Heat Equation

The four-parameter field $\mathcal{Z}^W((s, y), (t, x))$ solves the 1d SHE

$$\begin{cases} \partial_t \mathcal{Z}^W = \frac{1}{2} \Delta_x \mathcal{Z}^W + \hat{\beta} W \mathcal{Z}^W \\ \lim_{t \downarrow s} \mathcal{Z}^W((s, y), (t, x)) = \delta(y - x) \end{cases}$$

Checked directly from Wiener chaos expansion ([mild solution](#))

It is known that solutions to the SHE satisfy the properties above

Alternative approach (to check, OK for pinning [C., Sun, Zygouras 2016])

- Prove continuity by Kolmogorov criterion, showing that

$$\frac{\mathcal{Z}^W((s, y), (t, x))}{g_{t-s}(x - y)} \text{ is continuous also for } t = s$$

- Use continuity to prove semigroup for all times
- Use continuity to deduce positivity for close times, then bootstrap to arbitrary times using semigroup

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A naive approach

Consider DPRE in $d = 1$ (random walk + disorder)

$$\mathbf{P}^{\omega}(S) \propto e^{\sum_{n=1}^N \beta \omega(n, S_n)} \mathbf{P}^{\text{ref}}(S)$$

Can we define its **continuum analogue** (BM + disorder)? Naively

$$\mathcal{P}^W(\mathrm{d}B) \propto e^{\int_0^1 \hat{\beta} W(t, B_t) \mathrm{d}t} \mathcal{P}^{\text{ref}}(\mathrm{d}B)$$

\mathcal{P}^{ref} = law of BM $W(t, x)$ = white noise on \mathbb{R}^2 (space-time)

► $\int_0^1 W(t, B_t) \mathrm{d}t$ ill-defined. Regularization?

NO! The problem is more subtle (and interesting!)

Partition functions and f.d.d.

Start from **discrete**: distribution of DPRE at two times $0 < t < t' < 1$

$$\mathbf{P}_\delta^\omega(S_t^\delta = x, S_{t'}^\delta = x') = \frac{\mathbf{Z}_\delta^\omega((0,0), (t,x)) \mathbf{Z}_\delta^\omega((t,x), (t',x')) \mathbf{Z}_\delta^\omega((t',x'), (1,\star))}{\mathbf{Z}_\delta^\omega((0,0), (1,\star))}$$

(**drawing!**) Analogous formula for any finite number of times

Idea: Replace $\mathbf{Z}_\delta^\omega \rightsquigarrow \mathcal{Z}^\omega$ to *define* the law of continuum DPRE

Recall: to define a process $(X_t)_{t \in [0,1]}$ it is enough (Kolmogorov) to assign **finite-dimensional distributions** (f.d.d.)

$$\mu_{t_1, \dots, t_k}(A_1, \dots, A_k) = \mathbb{P}(X_{t_1} \in A_1, \dots, X_{t_k} \in A_k)$$

that are **consistent**

$$\mu_{t_1, \dots, t_j, \dots, t_k}(A_1, \dots, \mathbb{R}, \dots, A_k) = \mu_{t_1, \dots, t_{j-1}, t_{j+1}, \dots, t_k}(A_1, \dots, A_{j-1}, A_{j+1}, \dots, A_k)$$

The continuum 1d DPRE

- ▶ Fix $\hat{\beta} \in (0, \infty)$ (on which \mathcal{Z}^W depend) [recall that $\beta \sim \hat{\beta} \delta^{\frac{1}{4}}$]
- ▶ Fix space-time white noise W on $[0, 1] \times \mathbb{R}$ and a realization of continuum partition functions \mathcal{Z}^W satisfying the key properties (continuity, strict positivity, semigroup)

The **Continuum DPRE** is the process $(X_t)_{t \in [0,1]}$ with f.d.d.

$$\begin{aligned} & \frac{\mathcal{P}^W(X_t \in dx, X_{t'} \in dx')}{dx dx'} \\ & := \frac{\mathcal{Z}^W((0,0), (t,x)) \mathcal{Z}^W((t,x), (t',x')) \mathcal{Z}^W((t',x'), (1,\star))}{\mathcal{Z}^W((0,0), (1,\star))} \end{aligned}$$

- ▶ Well-defined by **strict positivity** of \mathcal{Z}^W
- ▶ Consistent by **semigroup** property

Relation with Wiener measure

The law of the continuum DPRE is a **random** probability

$$\mathcal{P}^W(X \in \cdot) \quad (\text{quenched law})$$

for the process $X = (X_t)_{t \in [0,1]}$ [Probab. kernel $\mathcal{S}'(\mathbb{R}) \rightarrow \mathbb{R}^{[0,1]}$]

Define a new law $\tilde{\mathbb{P}}$ (mutually absolutely continuous) for disorder **W** by

$$\frac{d\tilde{\mathbb{P}}}{d\mathbb{P}}(\mathbf{W}) = \mathcal{Z}^W((0,0), (1, \star))$$

Key Lemma

$$\mathcal{P}^{\text{ann}}(X \in \cdot) := \int_{\mathcal{S}'(\mathbb{R})} \mathcal{P}^W(X \in \cdot) \tilde{\mathbb{P}}(d\mathbf{W}) = \mathbb{P}(BM \in \cdot)$$

Proof. The factor \mathcal{Z}^W in $\tilde{\mathbb{P}}$ cancels the denominator in the f.d.d. for \mathcal{P}^W

Since $\mathbb{E}[\mathcal{Z}^W((s,y), (t,x))] = g_{t-s}(x-y)$ one gets f.d.d. of BM □

Absolute continuity properties

Theorem

$$\forall A \subseteq \mathbb{R}^{[0,1]} : P(BM \in A) = 1 \quad \Rightarrow \quad \mathcal{P}^W(X \in A) = 1 \quad \text{for } \mathbb{P}\text{-a.e. } W$$

Any given a.s. property of BM is an a.s. property of continuum DPRE, for a.e. realization of the disorder W

Corollary

$$\mathcal{P}^W(X \text{ has Hölder paths with exp. } \tfrac{1}{2}-) = 1 \quad \text{for } \mathbb{P}\text{-a.e. } W$$

We can thus realize \mathcal{P}^W as a law on $C([0,1], \mathbb{R})$, for \mathbb{P} -a.e. W

(More precisely: \mathcal{P}^W admits a modification with Hölder paths)

One is tempted to conclude that \mathcal{P}^W is absolutely continuous w.r.t. Wiener measure, for \mathbb{P} -a.e. W ...

NO! “ $\forall A$ ” and “for \mathbb{P} -a.e. W ” cannot be exchanged!

Singularity properties

Theorem

The law \mathcal{P}^W is **singular** w.r.t. Wiener measure, for \mathbb{P} -a.e. W .

$$\begin{array}{ll} \text{for } \mathbb{P}\text{-a.e. } W & \exists A = A_W \subseteq C([0, 1], \mathbb{R}) : \\ \mathcal{P}^W(X \in A) = 1 & \text{vs.} \quad P(BM \in A) = 0 \end{array}$$

Unlike discrete DPRE, there is **no continuum Hamiltonian**

$$\mathcal{P}^W(X \in \cdot) \not\propto e^{\mathcal{H}^W(\cdot)} P(BM \in \cdot)$$

Absolute continuity is lost in the scaling limit

In a sense, the laws \mathcal{P}^W are just *barely* not absolutely continuous w.r.t. Wiener measure (“stochastically absolutely continuous”)

Proof of singularity

Let $(X_t)_{t \in [0,1]}$ be the canonical process on $C([0,1], \mathbb{R})$ [$X_t(f) = f(t)$]

Let $\mathcal{F}_n := \sigma(X_{t_i^n} : t_i^n = \frac{i}{2^n}, 0 \leq i \leq 2^n)$ be the dyadic filtration

Fix a typical realization of W . Setting $\mathcal{P}^{\text{ref}} = \text{Wiener measure}$

$$R_n^W(X) := \frac{d\mathcal{P}^W|_{\mathcal{F}_n}}{d\mathcal{P}^{\text{ref}}|_{\mathcal{F}_n}}(X)$$

The process $(R_n^W)_{n \in \mathbb{N}}$ is a **martingale** w.r.t. \mathcal{P}^{ref} (**exercise!**)

Since $R_n^W \geq 0$, the martingale converges: $R_n^W \xrightarrow[n \rightarrow \infty]{\text{a.s.}} R_\infty^W$

- ▶ $\mathcal{P}^W \ll \mathcal{P}^{\text{ref}}$ if and only if $\mathcal{E}^{\text{ref}}[R_\infty^W] = 1$ (the martingale is UI)
- ▶ \mathcal{P}^W is **singular** w.r.t. \mathcal{P}^{ref} if and only if $R_\infty^W = 0$

Proof of singularity

It suffices to show that $R_n^W(X) \xrightarrow[n \rightarrow \infty]{} 0$ in $\mathbb{P} \otimes \mathcal{P}^{\text{ref}}$ -probability

Fractional moment

For \mathcal{P}^{ref} -a.e. X $\mathbb{E} \tilde{\mathbb{E}}[(R_n^W(X))^\gamma] \xrightarrow[n \rightarrow \infty]{} 0$ for some $\gamma \in (0, 1)$

$$R_n^W(X) = \frac{1}{\mathcal{Z}^W((0, 0), (1, \star))} \prod_{i=0}^{2^n-1} \frac{\mathcal{Z}^W((t_i^n, X_{t_i^n}), (t_{i+1}^n, X_{t_{i+1}^n}))}{g_{\frac{1}{2^n}}(X_{t_{i+1}^n} - X_{t_i^n})}$$

- Switch from \mathbb{E} to equivalent law $\tilde{\mathbb{E}}$ to cancel the denominator
- For fixed X , the $\mathcal{Z}^W((t_i^n, X_{t_i^n}), (t_{i+1}^n, X_{t_{i+1}^n}))$'s are independent

We need to exploit translation and scale invariance of their laws

Proof of singularity

Lemma 1 (Translation and scale invariance)

If we set $\Delta_i^n := \frac{X_{t_{i+1}^n} - X_{t_i^n}}{\sqrt{t_{i+1}^n - t_i^n}}$ we have

$$\frac{\mathcal{Z}_{\hat{\beta}}^W((t_i^n, X_{t_i^n}), (t_{i+1}^n, X_{t_{i+1}^n}))}{g_{\frac{1}{2^n}}(X_{t_{i+1}^n} - X_{t_i^n})} \stackrel{d}{=} \frac{\mathcal{Z}_{\frac{\hat{\beta}}{2^{n/4}}}^W((0, 0), (1, \Delta_i^n))}{g_1(\Delta_i^n)}$$

Lemma 2 (Expansion)

For $z \in \mathbb{R}$ and $\varepsilon \in [0, 1]$ (say)

$$\frac{\mathcal{Z}_{\varepsilon}^W((0, 0), (1, z))}{g_1(z)} = 1 + \varepsilon \mathbf{X}_z + \varepsilon^2 \mathbf{Y}_{\varepsilon, z}$$

$$\mathbb{E}[\mathbf{X}_z] = 0 \quad \mathbb{E}[\mathbf{X}_{\varepsilon, z}] = 0 \quad \mathbb{E}[\mathbf{X}_z^2] \leq C \quad \mathbb{E}[\mathbf{Y}_{\varepsilon, z}^2] \leq C \quad \text{unif. in } \varepsilon, z$$

Proof of singularity

By Taylor expansion, for fixed $\gamma \in (0, 1)$

$$\begin{aligned} \mathbb{E} \left[\left(\frac{\mathcal{Z}_\varepsilon^W((0,0), (1,z))}{g_1(z)} \right)^\gamma \right] &= \mathbb{E} \left[(1 + \varepsilon X_z + \varepsilon^2 Y_{\varepsilon,z})^\gamma \right] \\ &= 1 + \gamma \{ \varepsilon \mathbb{E}[X_z] + \varepsilon^2 \mathbb{E}[Y_{\varepsilon,z}] \} + \frac{\gamma(\gamma-1)}{2} \{ \varepsilon^2 \mathbb{E}[(X_x)^2] + \dots \} + \dots \\ &= 1 - c \varepsilon^2 \leq e^{-c \varepsilon^2} \end{aligned}$$

(★) First order terms vanish (★) $\gamma(\gamma-1) < 0$ (★) For some $c > 0$

Estimate is uniform over $z \in \mathbb{R} \rightsquigarrow$ We can set $z = \Delta_i^n$ and $\varepsilon = \frac{1}{2^{n/4}}$

$$\tilde{\mathbb{E}}[(R_n^W(X))^\gamma] = \prod_{i=0}^{2^n-1} \mathbb{E} \left[\left(\frac{\mathcal{Z}_\varepsilon^W((0,0), (1, \Delta_i^n))}{g_1(\Delta_i^n)} \right)^\gamma \right] \leq e^{-c \varepsilon^2 2^n} = e^{-c 2^{n/2}}$$

which vanishes as $n \rightarrow \infty$



Proof of Lemma 1

Introducing the dependence on $\hat{\beta}$

$$\begin{aligned} \mathcal{Z}_{\hat{\beta}}^W((s, y), (t, x)) &\stackrel{d}{=} \mathcal{Z}_{\hat{\beta}}^W((0, 0), (t - s, x - y)) \\ \mathcal{Z}_{\hat{\beta}}^W((0, 0), (t, x)) &\stackrel{d}{=} \frac{1}{\sqrt{t}} \mathcal{Z}_{\hat{\beta} t^{\frac{1}{4}}}^W\left((0, 0), \left(1, \frac{x}{\sqrt{t}}\right)\right) \end{aligned}$$

transl. invariance + diffusive rescaling (prefactor, new $\hat{\beta}$) (drawing!)

$$\begin{aligned} \mathcal{Z}^W((0, 0), (t, x)) &= g_t(x) + \hat{\beta} \int_{[0, t] \times \mathbb{R}} g_s(z) g_{t-s}(x - z) W(ds dz) + \dots \\ &= \frac{1}{\sqrt{t}} g_1\left(\frac{x}{\sqrt{t}}\right) + \frac{1}{\sqrt{t}} \left(\frac{\hat{\beta} t^{\frac{3}{4}}}{\sqrt{t}} \right) \int_{[0, t] \times \mathbb{R}} g_{\frac{s}{t}}\left(\frac{z}{\sqrt{t}}\right) g_{1-\frac{s}{t}}\left(\frac{x-z}{\sqrt{t}}\right) \frac{W(ds dz)}{t^{\frac{3}{4}}} + \dots \\ &= \text{OK!} \quad \square \end{aligned}$$

Convergence of discrete DPRE

- ▶ \mathbf{P}_δ^ω = law of **discrete** DPRE (recall that $S_t^\delta := \sqrt{\delta} S_{t/\delta}$)
“Rescaled RW S^δ moving in an i.i.d. environment ω ”
- ▶ \mathcal{P}^W = law of **continuum** DPRE
“BM moving in a white noise environment W ”

Both \mathbf{P}_δ^ω and \mathcal{P}^W are **random** probability laws on $E := C([0, 1], \mathbb{R})$ i.e. RVs (defined on different probab. spaces) taking values in $\mathcal{M}_1(E)$

Does \mathbf{P}_δ^ω converge in distribution toward \mathcal{P}^W as $\delta \rightarrow 0$?

$$\forall \psi \in C_b(\mathcal{M}_1(E) \rightarrow \mathbb{R}) : \quad \mathbb{E}[\psi(\mathbf{P}_\delta^\omega)] \xrightarrow[\delta \rightarrow 0]{} \mathbb{E}[\psi(\mathcal{P}^W)]$$

The answer is positive... almost surely ;-)

Statement for Pinning model proved in [C., Sun, Zygouras 2016]

Details need to be checked for DPRE (stronger assumptions on RW ?)

Universality

The convergence of \mathbf{P}_δ^ω toward \mathcal{P}^W is an instance of **universality**

There are **many** discrete DPRE:

- ▶ any RW S (zero mean, finite variance + technical assumptions)
- ▶ any (i.i.d.) disorder ω (finite exponential moments)

In the continuum ($\delta \rightarrow 0$) and weak disorder ($\beta \rightarrow 0$) regime, all these microscopic models \mathbf{P}_δ^ω give rise to a **unique** macroscopic model \mathcal{P}^W

Tomorrow we will see how the continuum model \mathcal{P}^W can tell quantitative information on discrete models \mathbf{P}_δ^ω (**free energy estimates**)

Convergence

How to prove convergence in distribution $\mathbf{P}_\delta^\omega \xrightarrow[\delta \rightarrow 0]{d} \mathcal{P}^W$?

Prove a.s. convergence through a suitable coupling of (ω, W)

Assume we have **convergence in distribution** of discrete partition functions to continuum ones, in the space of **continuum functions** of $(s, y), (t, x)$

$$\mathbf{Z}_\delta^\omega((s, y), (t, x)) \xrightarrow[\delta \rightarrow 0]{d} \mathcal{Z}^W((s, y), (t, x))$$

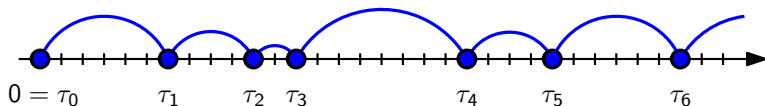
By Skorokhod representation theorem, there is a coupling of (ω, W) under which this convergence holds **a.s.**

Fix such a coupling: for a.e. (ω, W) the f.d.d. of \mathbf{P}_δ^ω converge weakly to those of \mathcal{P}^W . It only remains to prove tightness of $\mathbf{P}_\delta^\omega(\cdot)$.

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Ingredients: renewal process & disorder



Discrete **renewal process** $\tau = \{0 = \tau_0 < \tau_1 < \tau_2 < \dots\} \subseteq \mathbb{N}_0$

Gaps $(\tau_{i+1} - \tau_i)_{i \geq 0}$ are **i.i.d.** with polynomial-tail distribution:

$$\mathbf{P}^{\text{ref}}(\tau_1 = n) \sim \frac{c_K}{n^{1+\alpha}}, \quad c_K > 0, \quad \alpha \in (0, 1)$$

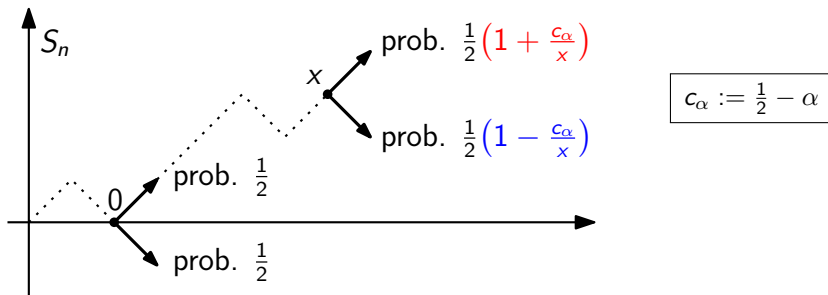
$\tau = \{n \in \mathbb{N}_0 : S_n = 0\}$ zero level set of a Markov chain $S = (S_n)_{n \geq 0}$

Disorder $\omega = (\omega_i)_{i \in \mathbb{N}}$: i.i.d. real random variables with law \mathbb{P}

$$\lambda(\beta) := \log \mathbb{E}[e^{\beta \omega_1}] < \infty \quad \mathbb{E}[\omega_1] = 0 \quad \text{Var}[\omega_1] = 1$$

Bessel random walks

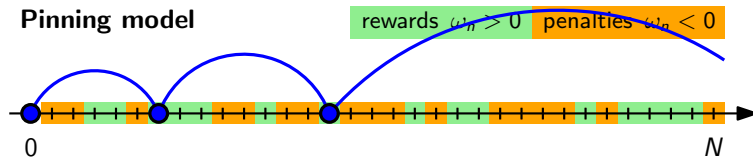
For $\alpha \in (0, 1)$ the α -Bessel random walk is defined as follows:



- ▶ $(\alpha = \frac{1}{2})$ **no drift** ($c_\alpha = 0$) \rightsquigarrow simple random walk
- ▶ $(\alpha < \frac{1}{2})$ drift **away** from the origin ($c_\alpha > 0$)
- ▶ $(\alpha > \frac{1}{2})$ drift **toward** the origin ($c_\alpha < 0$)

Disordered pinning model

Pinning model



$N \in \mathbb{N}$ (system size) $\beta \geq 0$, $h \in \mathbb{R}$ (disorder strength, bias)

The pinning model

Gibbs change of measure $\mathbf{P}_N^\omega = \mathbf{P}_{N,\beta,h}^\omega$ of the renewal distribution \mathbf{P}^{ref}

$$\frac{d\mathbf{P}_N^\omega}{d\mathbf{P}^{\text{ref}}}(\tau) := \frac{1}{Z_N^\omega} \exp \left(\sum_{n=1}^N (\beta \omega_n + h - \lambda(\beta)) \mathbb{1}_{\{n \in \tau\}} \mathbb{1}_{\{S_n=0\}} \right)$$

The phase transition

How are the typical paths τ of the pinning model \mathbf{P}_N^ω ?

Contact number $\mathcal{C}_N := |\tau \cap (0, N]| = \sum_{n=1}^N \mathbb{1}_{\{n \in \tau\}} = \sum_{n=1}^N \mathbb{1}_{\{S_n=0\}}$

Theorem (phase transition)

\exists *continuous, non decreasing, deterministic critical curve* $h_c(\beta)$:

► *Localized regime*: for $h > h_c(\beta)$ one has $\mathcal{C}_N \approx N$

$$\exists \mu = \mu_{\beta, h} > 0 : \quad \mathbf{P}_N^\omega \left(\left| \frac{\mathcal{C}_N}{N} - \mu \right| > \varepsilon \right) \xrightarrow{N \rightarrow \infty} 0 \quad \omega\text{-a.s.}$$

► *Deocalized regime*: for $h < h_c(\beta)$ one has $\mathcal{C}_N = O(\log N)$

$$\exists A = A_{\beta, h} > 0 : \quad \mathbf{P}_N^\omega \left(\frac{\mathcal{C}_N}{\log N} > A \right) \xrightarrow{N \rightarrow \infty} 0 \quad \omega\text{-a.s.}$$

Estimates on the critical curve

For $\beta = 0$ (homogeneous pinning, no disorder) one has $h_c(0) = 0$

What is the behavior of $h_c(\beta)$ for $\beta > 0$ small ?

Theorem ($P(\tau_1 = n) \sim \frac{c_K}{n^{1+\alpha}}$)

- ▶ ($\alpha < \frac{1}{2}$) *disorder is irrelevant*: $h_c(\beta) = 0$ for $\beta > 0$ small
[Alexander] [Toninelli] [Lacoin] [Cheliotis, den Hollander]
- ▶ ($\alpha \geq \frac{1}{2}$) *disorder is relevant*: $h_c(\beta) > 0$ for all $\beta > 0$
 - ▶ ($\alpha > 1$) $h_c(\beta) \sim C \beta^2$ with explicit $C = \frac{\alpha}{1+\alpha} \frac{1}{2E(\tau_1)}$
[Berger, C., Poisat, Sun, Zygouras]
 - ▶ ($\frac{1}{2} < \alpha < 1$) $C_1 \beta^{\frac{2\alpha}{2\alpha-1}} \leq h_c(\beta) \leq C_2 \beta^{\frac{2\alpha}{2\alpha-1}}$ $h_c(\beta) \sim \hat{C} \beta^{\frac{2\alpha}{2\alpha-1}}$
using continuum part. funct.!
[Derrida, Giacomin, Lacoin, Toninelli] [Alexander, Zygouras] [C., Torri, Toninelli]
 - ▶ ($\alpha = \frac{1}{2}$) $h_c(\beta) = e^{-\frac{c+o(1)}{\beta^2}}$ [Giacomin, Lacoin, Toninelli] [Berger, Lacoin]

Discrete free energy and critical curve

Partition function
$$\mathbf{Z}_N^\omega := \mathbb{E} \left[e^{H_N(\tau)} \right] = \mathbb{E} \left[e^{\sum_{n=1}^N (h + \beta \omega_n - \Lambda(\beta)) \mathbb{1}_{\{n \in \tau\}}} \right]$$

Consider first the regime of $N \rightarrow \infty$ with **fixed** β, h

► **Free energy**
$$F(\beta, h) := \lim_{N \rightarrow \infty} \frac{1}{N} \log \mathbf{Z}_N^\omega \geq 0 \quad \mathbb{P}(d\omega)\text{-a.s.}$$

$$\mathbf{Z}_N^\omega \geq \mathbb{E} \left[e^{H_N(\tau)} \mathbb{1}_{\{\tau \cap (0, N] = \emptyset\}} \right] = \mathbb{P}(\tau \cap (0, N] = \emptyset) \sim \frac{(\text{const.})}{N^\alpha}$$

► **Critical curve**
$$h_c(\beta) = \sup\{h \in \mathbb{R} : F(\beta, h) = 0\} \quad \text{non analyticity!}$$

(convexity)
$$\frac{\partial F(\beta, h)}{\partial h} = \lim_{N \rightarrow \infty} \mathbf{E}_N^\omega \left[\frac{\mathcal{C}_N}{N} \right] \begin{cases} > 0 & \text{if } h > h_c(\beta) \\ = 0 & \text{if } h < h_c(\beta) \end{cases}$$

$F(\beta, h)$ and $h_c(\beta)$ depend on the law of τ and ω

Universality as $\beta, h \rightarrow 0$? **YES**, connected to continuum model

Continuum partition functions

Build **continuum partition functions** for Pinning Model with $\alpha \in (\frac{1}{2}, 1)$ (disorder relevant) following “usual” approach [C, Sun, Zygouras 2015+]

We need to rescale

$$\beta = \beta_N = \frac{\hat{\beta}}{N^{\alpha-1/2}} \quad h = h_N = \frac{\hat{h}}{N^\alpha}$$

One has $\mathbf{Z}_N^\omega \xrightarrow[N \rightarrow \infty]{d} \mathcal{Z}^w$ with

$$\mathcal{Z}^w := 1 + C \int_{0 < t < 1} \frac{dW_t^{\hat{\beta}, \hat{h}}}{t^{1-\alpha}} + C^2 \int_{0 < t < t' < 1} \frac{dW_t^{\hat{\beta}, \hat{h}} dW_{t'}^{\hat{\beta}, \hat{h}}}{t^{1-\alpha} (t' - t)^{1-\alpha}} + \dots$$

where $W_t^{\hat{\beta}, \hat{h}} := \hat{\beta} W_t + \hat{h} t$ and $C = \frac{\alpha \sin(\alpha\pi)}{\pi c_K}$

Continuum free energy

In analogy with the discrete model, define

Continuum free energy $\mathcal{F}(\hat{\beta}, \hat{h}) := \lim_{t \rightarrow \infty} \frac{1}{t} \log \mathcal{Z}_{\hat{\beta}, \hat{h}}^W(0, t) \quad \text{a.s.}$

(existence and self-averaging need some work)

Again $\mathcal{F}(\hat{\beta}, \hat{h}) \geq 0$ and define

Continuum critical curve $\mathcal{H}_c(\hat{\beta}) := \sup \{ \hat{h} \in \mathbb{R} : \mathcal{F}(\hat{\beta}, \hat{h}) = 0 \}$

Scaling relations

$$\forall c > 0 : \quad \mathcal{Z}_{\hat{\beta}, \hat{h}}^W(ct) \stackrel{d}{=} \mathcal{Z}_{c^{\alpha - \frac{1}{2}}\hat{\beta}, c^{\alpha}\hat{h}}^W(t) \quad (\text{Wiener chaos exp.})$$

$$\mathcal{F}(c^{\alpha - \frac{1}{2}}\hat{\beta}, c^{\alpha}\hat{h}) = c \mathcal{F}(\hat{\beta}, \hat{h})$$

$$\mathcal{H}_c(\hat{\beta}) = \mathcal{H}_c(1) \hat{\beta}^{\frac{2\alpha}{2\alpha-1}}$$

Interchanging the limits

Can we relate continuum free energy to the discrete one?

By construction of continuum partition functions

$$\mathcal{Z}_{\hat{\beta}, \hat{h}}^W(t) \stackrel{d}{=} \lim_{N \rightarrow \infty} \mathbf{Z}_{\beta_N, h_N}^{\omega}(Nt)$$

Assuming uniform integrability of $\log \mathbf{Z}^{\omega}$ (OK)

$$\mathcal{F}(\hat{\beta}, \hat{h}) = \lim_{t \rightarrow \infty} \frac{1}{t} \mathbb{E}[\log \mathcal{Z}_{\hat{\beta}, \hat{h}}^W(t)] = \lim_{t \rightarrow \infty} \frac{1}{t} \lim_{N \rightarrow \infty} \mathbb{E}[\log \mathbf{Z}_{\beta_N, h_N}^{\omega}(Nt)]$$

Assuming we can interchange the limits $N \rightarrow \infty$ and $t \rightarrow \infty$

$$\mathcal{F}(\hat{\beta}, \hat{h}) = \lim_{N \rightarrow \infty} N \lim_{t \rightarrow \infty} \frac{1}{Nt} \mathbb{E}[\log \mathbf{Z}_{\beta_N, h_N}^{\omega}(Nt)] = \lim_{N \rightarrow \infty} N F(\beta_N, h_N)$$

Setting $\delta = \frac{1}{N}$ for clarity, we arrive at...

Interchanging the limits

Conjecture

$$\mathcal{F}(\hat{\beta}, \hat{h}) = \lim_{\delta \rightarrow 0} \frac{F(\hat{\beta}\delta^{\alpha-\frac{1}{2}}, \hat{h}\delta^{\alpha})}{\delta}$$

Theorem [C., Toninelli, Torri 2015]

For all $\hat{\beta} > 0$, $\hat{h} \in \mathbb{R}$ and $\eta > 0$ there is $\delta_0 > 0$ such that $\forall \delta < \delta_0$

$$\mathcal{F}(\hat{\beta}, \hat{h} - \eta) \leq \frac{F(\hat{\beta}\delta^{\alpha-\frac{1}{2}}, \hat{h}\delta^{\alpha})}{\delta} \leq \mathcal{F}(\hat{\beta}, \hat{h} + \eta)$$

This implies Conj. and

$$h_c(\beta) \sim \mathcal{H}_c(\beta) \sim \mathcal{H}_c(1) \beta^{\frac{2\alpha}{2\alpha-1}}$$

For **any** discrete Pinning model with $\alpha \in (\frac{1}{2}, 1)$, the free energy $\mathbf{F}(\beta, h)$ and the critical curve $h_c(\beta)$ have a **universal shape** in the regime $\beta, h \rightarrow 0$

Interchanging the limits

Very delicate result. How to prove it?

- ▶ Assume that there is a continuum Hamiltonian:

$$\mathbf{Z}^\omega = \mathbf{E}[e^{\mathbf{H}_{Nt}^\omega}] \quad \mathcal{Z}^W = \mathcal{E}[e^{\mathcal{H}_t^W}]$$

- ▶ Couple \mathbf{H}_{Nt}^ω and \mathcal{H}_t^W on the same probability space in such a way that the difference $\Delta_{N,t} := \mathbf{H}_{Nt}^\omega - \mathcal{H}_t^W$ is “small”
- ▶ Deduce that

$$\mathbb{E}[\log \mathbf{Z}^\omega] \leq \mathbb{E}[\log \mathcal{Z}^W] + \log \mathbb{E}[e^{\Delta_{N,t}}]$$

and show that the last term is “negligible”

Problem: there is no continuum Hamiltonian!

Solution: perform **coarse-graining** and define an “effective” Hamiltonian

The DPRE case

What about the DPRE?

We can still define discrete $F(\beta)$ and continuum $\mathcal{F}(\hat{\beta})$ free energy

Since $\mathcal{F}(\hat{\beta}) \sim \mathcal{F}(1)\beta^4$ we can hope that

$$F(\beta) \sim \mathcal{F}(1)\beta^4 \quad \text{as } \beta \rightarrow 0$$

provided the “interchanging of limits” is justified

N. Torri is currently working on this problem. A finer coarse-graining is needed, together with sharper estimates on continuum partition functions

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